

# A second order equation for Schrödinger Bridges and applications

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# Moving towards an “unexpected configuration”

Back in 1932, Schrödinger described the following experiment

- At time  $t = 0$  consider  $N$  **Brownian particles on the sphere**  $(X_t^i)_{t \leq 1, i \leq N}$  arranged according to  $\mu$ ,

$$\mu = \frac{1}{N} \sum_{i=1}^N \delta_{X_0^i}.$$

- Let each particle move independently for a unit of time and look at their configuration  $\nu$  at time 1,

$$\nu = \frac{1}{N} \sum_{i=1}^N \delta_{X_1^i}$$

- Assume that  $\nu$  is an **unexpected configuration**, i.e.  $\nu \neq \text{vol}$ .

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# Schrödinger problem: Large Deviations formulation

- $\mathbf{P}$  = stationary process for generator  $\mathcal{L}$

$$\mathcal{L} = \frac{1}{2}\Delta - \nabla U \cdot \nabla$$

- $\mu, \nu$  initial and final configurations
- $\mathcal{H}(\cdot|\mathbf{P})$  = Relative entropy on path space

## Schrödinger problem (SP)

$$\min \mathcal{H}(\mathbf{Q}|\mathbf{P}), \quad X_0 \# \mathbf{Q} = \mu, X_1 \# \mathbf{Q} = \nu$$

- The optimal value  $\mathcal{T}_{\mathcal{H}}(\mu, \nu)$  of the Schrödinger problem is called **entropic transportation cost**.
- The **optimal solution** of SP is the **Schrödinger bridge (SB)** between  $\mu$  and  $\nu$ .

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The SB models a particle system that

- has to **minimize entropy**  $\Rightarrow$  particles are willing to arrange themselves according to the equilibrium configuration **m**.
- has to reach an **unexpected** final configuration, which looks very different from **m**

Thus we expect the dynamics to be divided into two phases.

- 1 Entropy minimization dominates: SB relaxes to **m**
- 2 the influence of the final configuration prevails: particles start arranging according to  $\nu$  and drifts away from **m**.

At any given time  $t$ , how far is  $\mu_t$  from **m** ?

Goal of the talk: answer this question.

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# Finding an equation for the Schrödinger bridge

We call  $(\mu_t)$  the **marginal flow** of SB, and  $(v_t)$  its “speed”

- Quickest way to minimize entropy: **gradient flow of the entropy**

$$v_t = -\nabla \mathcal{H}(\mu_t)$$

- But if particles go along the gradient flow they don't reach the “unexpected configuration”  $\nu$
- **Cheap trick**: differentiate in time the gradient flows

$$\frac{D}{dt} v_t = -\nabla^2 \mathcal{H}(\mu_t) v_t = \nabla^2 \mathcal{H}(\mu_t) \nabla \mathcal{H}(\mu_t)$$

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# The Riemannian manifold of Optimal transport

- To make sense of the calculations above we need to choose some **metric structure** on the space of probability measures
- The theory of **optimal transport** can be used to construct one such structure.
- The **continuity equation** gives the notion of speed for a curve

$$\partial_t \mu + \nabla \cdot (\mu_t v_t) = 0$$

- The **convective derivative** can be used to make sense of the **covariant derivative**.

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# A second order equation

Our heuristic can be rigorously justified within this structure!

Theorem (C.' 17)

*The marginal flow  $(\mu_t)$  of the Schrödinger bridge solves*

$$\frac{\mathbf{D}}{dt} v_t = \frac{1}{8} \nabla^{\mathcal{W}} \mathcal{I}_U(\mu_t)$$

where

- $\frac{\mathbf{D}}{dt}$  is the **covariant derivative**
- $v_t$  is the **velocity field** of  $(\mu_t)$
- $\mathcal{I}_U$  is the **modified Fisher information**.
- $\nabla^{\mathcal{W}}$  is the **gradient in the OT sense**.



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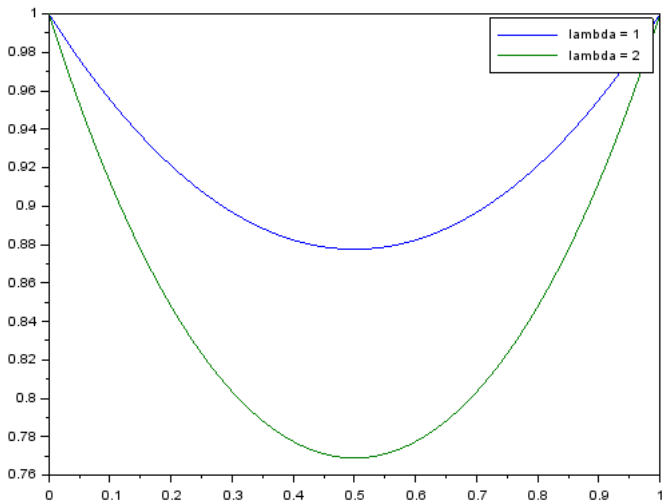
Let  $M$  be compact and  $(\mu_t)$  be the entropic interpolation between  $\mu$  and  $\nu$ . The Bakry Émery condition

$$\mathfrak{Ric} + \nabla \nabla U \geq \lambda$$

implies the following bound for all  $\mu, \nu \in \mathcal{P}(M)$

$$\mathcal{H}_U(\mu_t) \leq \frac{1 - \exp(-\alpha(1-t))}{1 - \exp(-\alpha)} \mathcal{H}_U(\mu) + \frac{1 - \exp(-\alpha t)}{1 - \exp(-\alpha)} \mathcal{H}_U(\nu) - \frac{\cosh(\frac{\alpha}{2}) - \cosh(\alpha(t - \frac{1}{2}))}{\sinh(\frac{\alpha}{2})} \mathcal{T}_{\mathcal{H}}(\mu, \nu)$$

# Evolution of the marginal entropy





## Theorem (C.'17)

Let  $m$  be the invariant measure associated with  $\mathcal{L}$ . Then

$$\mathcal{T}_{\mathcal{H}_U}(\mu, \mathbf{m}) \leq \frac{1}{1 - \exp(-\frac{\alpha}{2})} \mathcal{H}_U(\mu)$$

- A priori bound on the entropic transportation cost
- The entropic cost grows at most linearly with the marginal entropies
- Upper bound for the distance between the Schrödinger bridge and the target law  $\mathbf{P}$
- Partial converse to the fact that joint entropies dominate marginal entropies

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If we consider the SB constructed with

$$\mathcal{L}^\varepsilon = \frac{\varepsilon}{2} \Delta$$

Theorem ( Mikami '04, Léonard '12 )

*The following limits hold as  $\varepsilon \downarrow 0$ .*

*Schrödinger bridge  $\rightarrow$  displacement interpolation*

*and*

$$\varepsilon \mathcal{T}_{\mathcal{H}_U}^\varepsilon(\mu, \nu) \rightarrow \frac{1}{2} W_2^2(\mu, \nu)$$

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## Reciprocal characteristic

For a given potential  $U$ , it is the vector field

$$\nabla \mathcal{U}, \quad \text{where } \mathcal{U} = \frac{1}{2} |\nabla U|^2 - \frac{1}{2} \Delta U$$

In the flat case  $M = \mathbb{R}^d$  we obtained the following

## Theorem (C.'17)

*Let  $(\mu_t)$  be the marginal flow of (SB). If  $\mathcal{U}$  is convex and  $\mu_t$  is a log-concave law for all  $t$ , then*

$$t \mapsto \mathcal{I}_U(\mu_t) \quad \text{is convex}$$

- Convexity points of  $\mathcal{I}_U = \log$  concave measures?



# Grazie!

beginframeSchrödinger '33

*“ Imaginez que vous observez un **système de particules en diffusion**, qui soient en équilibre thermodynamique. Admettons qu' à un instant donné 0 vous les ayez trouvées en **répartition à peu près uniforme** et que à 1 vous ayez trouvé un **écart spontané et considérable par rapport à cette uniformité**. On vous demande de quelle manière cet écart sest produit. Quelle en est la **manière la plus probable** ? ”*



G. Conforti.

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