Local Laws in RMT

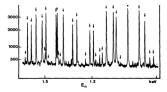
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June 2, 2015

Wigner's Nuclear Physics

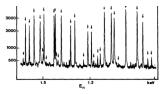
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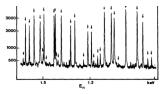
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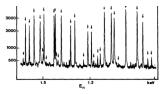
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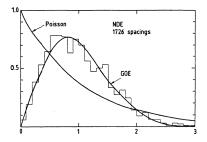
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Question: How to understand the spacing between peaks in the above data?

- Structure of large nuclei is very complicated
- Not feasible to compute these energies either theoretically or numerically
- Wigner's idea: the Hamiltonian can be modeled by a very large random matrix

Data for spacing between energies with eigenvalue spacing overlaid:



Definition

Let *h* be a probability distribution on \mathbb{C} . A **Wigner matrix** is a $N \times N$ matrix that is:

- Hermitian,
- entries are iid complex random variables with distribution h,
- mean 0 and variance 1/N.

The variance is chosen so the spectrum has compact support.

Universality

Consider a Wigner matrix (Hermitian $N \times N$ matrix with iid entries drawn from any distribution h). Let

- $\mathcal{N}[a; b] =$ number of eigenvalues in the interval [a; b]
- ρ_{sc} is a probability density on \mathbb{R}

$$\rho_{sc}(E) = \begin{cases} \frac{1}{2\pi}\sqrt{1 - \frac{E^2}{4}}, & \text{if } |E| \le 2\\ 0 & \text{if } |E| > 2 \end{cases}$$

Wigner's key result – universality: For $\delta > 0$ and all a, b, Wigner proved that

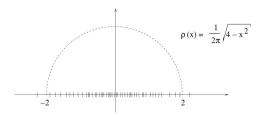
$$\lim_{N\to\infty} \mathbb{P}\left(\left| \frac{\mathcal{N}[a;b]}{\mathcal{N}(b-a)} - \frac{1}{(b-a)} \int_{a}^{b} \rho_{sc}(s) ds \right| \geq \delta \right) = 0$$

Universality - the big picture

- Want to derive macroscopic properties of large systems with unknown or random interactions of their constituents. Examples:
 - laws of thermodynamics
 - law of large numbers
- Random matrices model various physical systems
 - quantum and wave chaos
 - many particle systems
 - scattering energies of large atoms
- For random matrices, universality can be defined and proved
 - A lot done for Wigner matrices
 - Still many open problems for other ensembles, such as non-Hermitian matrices
 - Universality of fluctuations?

Universality in random matrices suggests a mechanism of universality in physical systems

Zooming in



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Counting number of eigenvalues of Wigner matrices in short intervals:

- In Wigner' work, interval [a, b] independent of N
- What if we take [a(N), b(N)] with $b(N) a(N) \rightarrow 0$ as $N \rightarrow \infty$?

 Erdös-Schlein-Yau 2008, 2009, 2010. Let δ > 0 and E away from spectral edge. There exists a constant K_δ such that with η = K_δ/N,

$$\mathbb{P}\left(\left|\frac{\mathcal{N}(\mathsf{E}-\eta,\mathsf{E}+\eta)}{2\mathsf{N}\eta}-\rho_{\mathsf{sc}}(\mathsf{E})\right|>\delta\right)<\mathsf{C}\mathsf{e}^{-c\delta^2\sqrt{\mathsf{N}\eta}}$$

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For any sequence $\eta(N) \rightarrow 0$

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Need: Gaussian decay, two derivatives, and $\int \left(\frac{f'(x)}{f(x)}\right)^2 f(x) dx < \infty$.

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$$\mathbb{P}\left(\left|\lambda_{\alpha}-\gamma_{\alpha}\right| \geq \frac{K \log N}{N} \left(\frac{N}{\hat{\alpha}}\right)^{\frac{1}{3}}\right) \leq \frac{(Cq)^{cq^{2}}}{K^{q}}$$
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for all $N > N_0, K > 0$, $q \in \mathbb{N}$ with $q < N^{\varepsilon}$.

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- Eigenvalues are expected to fluctuate on the scale $\sqrt{\log N}/N$ (Gustavson for GUE)
- Improved from $\frac{(\log N)^{\log N}}{N}$ (up to constants) in prior work of Erdös-Yau-Yin

2-pt correlation function:

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where p_N is the joint probability distribution of the eigenvalues;

$$\frac{1}{\rho(E)}p^{(2)}\left(E+\frac{x_1}{\rho(E)N},E+\frac{x_2}{\rho(E)N}\right) \to \det\left(1-\left(\frac{\sin(\pi x_1-\pi x_2)}{\pi x_1-\pi x_2}\right)^2\right)$$

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- Tao Vu used a moment method to prove it for Wigner

Other problems solved and unsolved

Local laws in other ensembles

- Covariance matrices mostly solved!
- Non-Hermitian iid matrices somewhat solved
- Sparse somewhat solved
- Heavy-tailed and heavier-tailed Wigner a little bit of work done but mostly open

Technicalities

$$G_{jj} = \frac{1}{h_{jj} - z - \mathbf{a}_j^* G^{(j)} \mathbf{a}_j}$$
(5.1)

Let $\Lambda = m - m_{sc}$, then

$$G_{jj} = -\frac{1}{m_{sc} + z + \Lambda + \Upsilon_j}$$
(5.2)

where

$$\Upsilon_j = -h_{jj} - \frac{1}{N}G_{jj} - \frac{1}{N}\sum_{k\neq j}(G_{kk} - G_{kk}^{(j)}) + (I - \mathbb{E}_j)\mathbf{a}_j^*G^{(j)}\mathbf{a}_j$$

where \mathbb{E}_i denotes the expectation with respect to \mathbf{a}_i .

Solve for the fluctuation

We set
$$R = \frac{1}{N} \sum_{j} \Upsilon_{j} G_{jj}$$
, then so that Eq. (5.2) reads

$$m_{sc}\Lambda^2 + (m_{sc}^2 - 1)\Lambda + m_{sc}R = 0.$$
 (5.3)

•

We solve the equation for Λ and choose the correct solution

$$\Lambda = \frac{-(m_{sc}^2 - 1) + \sqrt{(m_{sc}^2 - 1)^2 - 4m_{sc}^2 R}}{2m_{sc}}$$

$$|\Lambda| \le C \frac{|m_{sc}||R|}{\sqrt{|m_{sc}^2 - 1|^2 + 4|m_{sc}|^2|R|}} \le C \min\left\{\frac{|R|}{|m_{sc}^2 - 1|}, \sqrt{|R|}\right\}.$$
 (5.4)

Hanson & Wright: fluctuations of quadratic forms

For j = 1, ..., N let $x_j = \text{Re } x_j + i \text{ Im } x_j$, where {Re x_j , Im x_j } $_{j=1}^N$ is a sequence of 2N real iid random variables, whose common distribution ν has subgaussian decay. Let $A = (a_{ij})$ be a $N \times N$ complex matrix. Then there exist constants c, C > 0 such that, for any $\delta > 0$

$$\mathbb{P}\left(\left|\sum_{i,j=1}^{N} a_{ij}\left(x_{i}\bar{x}_{j} - \mathbb{E}x_{i}\bar{x}_{j}\right)\right| \geq \delta\right) \leq Ce^{-c\delta/\sqrt{\operatorname{Tr}A^{*}A}}.$$
(5.5)

Let $z = E + i\eta$. For any j = 1, ..., N, set

$$Z_j = (I - \mathbb{E}_j) \mathbf{a}_j^* G^{(j)} \mathbf{a}_j \,. \tag{5.6}$$

• Setting A = G and integrating over δ

$$\mathbb{E}|Z_j|^{2q} \leq (Cq)^{2q} \left(\frac{(\operatorname{Im} m_{sc})^q + \mathbb{E}|\Lambda|^q}{(N\eta)^q} + \frac{1}{(N\eta)^{2q}} \right), \quad (5.7)$$

• $\mathbb{E}|G_{12}|^{2q} < \frac{C_q}{(N\eta)^q}$ - tricky!

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